

# Gate Leakage vs. NBTI in Plasma Nitrided Oxides: Characterization, Physical Principles, and Optimization

\*<sup>1</sup>A.E. Islam, <sup>2</sup>G. Gupta, <sup>2</sup>S. Mahapatra, <sup>3</sup>A.T. Krishnan, <sup>4</sup>K. Ahmed, <sup>4</sup>F. Nouri, <sup>5</sup>A. Oates, <sup>1</sup>M.A. Alam

\*Email: aeislam@purdue.edu, Phone: 765-494-5988, Fax: 765-494-6441

<sup>1</sup>Electrical and Computer Engineering, Purdue University, West Lafayette, IN, USA; <sup>2</sup>Department of EE, IIT Bombay, India; <sup>3</sup>Texas Instruments, Dallas, TX, USA; <sup>4</sup>Applied Materials, Inc., Santa Clara, CA, USA; <sup>5</sup>Taiwan Semiconductor Manufacturing Corp., Hsin-Chu, Taiwan.

## Abstract

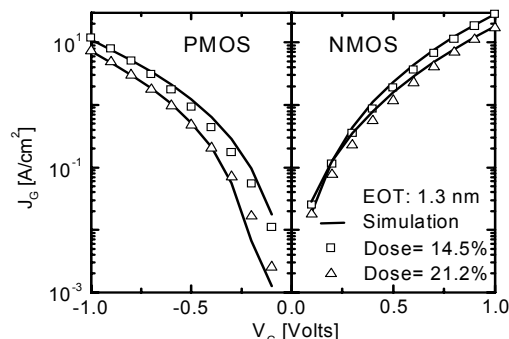
Since nitrided oxides improve gate leakage at the expense of NBTI, one must optimize nitrogen concentration in oxinitride samples for reliable performance and reduced power dissipation. Here, we analyze wide range of NBTI stress data to develop a predictive model for gate leakage and first self-consistent model for field acceleration within R-D framework. This model anticipates a novel design diagram for co-optimization of leakage and NBTI for arbitrary nitrogen concentration and effective oxide thickness.

## 1. Introduction

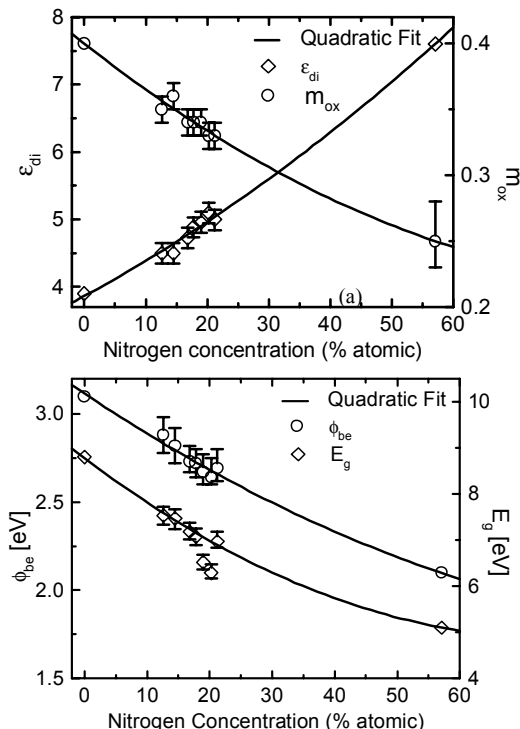
Since nitrided oxides (SiON) improve gate leakage ( $J_G$ ) [1-6] at the expense of NBTI performance ( $\Delta V_T$ ) [7,8], one must necessarily optimize  $N_2$  concentration (%N) in gate-oxides for high-performance ICs. Despite its importance, however, a quantitative analysis of leakage/NBTI trade-off (as a function of %N), has never been reported and the question “Is co-optimization of NBTI/leakage possible at any %N?” has never been answered. In this paper, we simultaneously measure gate leakage and delay-free NBTI over broad range of stress-fields, stress-temperatures and %N, model gate leakage current ( $J_G$ ) and NBTI degradation within a theoretically consistent framework (hole-assisted thermal generation of interface traps) of field-dependent R-D model, and conclude that although there is no optimum %N for NBTI/leakage, the reduction in  $J_G$  at NBTI-limited %N (~15-25%, depending on failure criterion) can be significant and would reduce power dissipation without affecting NBTI-margin.

## 2. Gate Leakage

Comprehensive simulation [9] (which includes the effects of multi-subband electron/hole quantization, poly-depletion, etc.) of the measured  $J_G$ - $V_G$  for both N- and PMOS (Fig. 1) was done to extract the model parameters as a function of %N (Fig. 2). We assume that any variation in the spatial-profile of nitrogen results only in second-order correction to calculated  $J_G$ . Contrary to popular belief [2-6, 10], the oxide parameters do not scale linearly with %N. All the parameters have approximately a quadratic fit with %N. Here, effective oxide thickness ( $EOT$ ) is obtained from simulation of CV, and physical thickness ( $T_{PHY}$ ) and %N are determined by XPS [11]. These %N-dependent parameters are used to calculate  $J_G(N, EOT)$  for arbitrary %N and  $EOT$ , as shown in Fig. 10b.



**Fig. 1:** Experimental leakage current ( $J_G$ ) vs. gate voltage ( $V_G$ ) curves for SiON devices (14.5% and 21.2%  $N_2$  dose) in inversion region, fitted using simulation.



**Fig. 2:** Variation of gate leakage model parameters with %N (0% for  $SiO_2$ , 57.1% for  $Si_3N_4$ ) (a) relative permittivity for SiON dielectric,  $\epsilon_{di}$  and oxide effective mass,  $m_{ox}$  (b) barrier height for electron,  $\phi_{be}$  and bandgap,  $E_g$ . For  $Si_3N_4$ , parameters are obtained from references [2-6] ( $\phi_{be}$ : 2.1 eV,  $E_g$ : 5.1 eV,  $\epsilon_{di}$ : 7.6,  $m_{ox}$ : 0.23~0.28). The error bars show the variation in parameters expected due to  $\pm 0.5$  Å error in measurement of  $T_{PHY}$  by XPS.

### 3. Reliability based on NBTI

NBTI performance for different SiON devices is monitored here by calculating safe operating voltage ( $V_{safe}$ ), which is defined as the stress voltage at which device can operate up to its lifetime ( $t_{life}$ ) without crossing the failure criteria (e.g. 60mV of  $\Delta V_T$  [1]). As shown in Fig. 3, degradation data taken at different stress voltages ( $V_1, V_2, V_3, V_4$ ) can be fitted with analytical expressions, based on the field-dependent R-D model to be discussed Section 3.1. Such fitting can be used to extract lifetimes ( $t_1, t_2, t_3, t_4$ ) at stress voltages and  $V_{safe}$ , based on failure criteria. The calculated  $V_{safe}(N, EOT)$  for arbitrary %N and  $EOT$  is plotted in Fig. 10a.

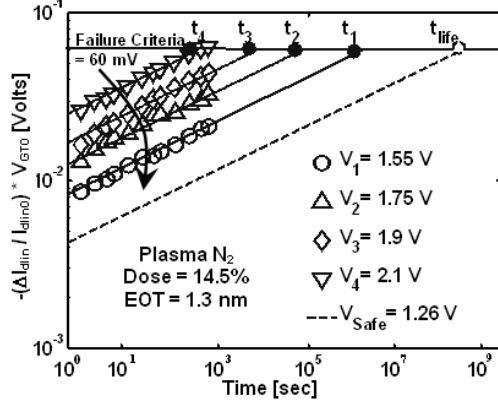


Fig. 3: Procedure for obtaining lifetimes ( $t_1, t_2, t_3, t_4$ ) at different stress voltages ( $V_1, V_2, V_3, V_4$ ) and  $V_{safe}$  using NBTI model, discussed in section 3.1. Proper  $V_G$ - $E_{ox}$  transformation is used here and  $I_{dlin0}$  was obtained within 1ms of application of stress.

#### 3.1 Field-Dependent R-D model

Since NBTI is field-dependent [12], the  $E_{ox}$  (calculated either from direct integration of CV, or approximated using proper  $V_G$ - $E_{ox}$  relationship) dependence of NBTI parameters needs to be established. The R-D model anticipates that [13]

$$\Delta N_{IT} \sim (k_{f0} N_0 / k_{r0})^{2/3} \{ D_0 \exp(-E_D / k_B T) t \}^{1/6} \quad (1)$$

for neutral- $H_2$  diffusion [13-15], which is supported by delay-free NBTI measurements [13-16]. Since the diffusion coefficient pre-factor,  $D_0 = l^2 \nu_0$  ( $l$  is lattice constant,  $\nu_0$  is attempt frequency) is field-independent, NBTI field-dependence must be encapsulated in the ratio of forward ( $k_{f0}$ ) and reverse ( $k_{r0}$ ) dissociation constants, i.e.  $k_{f0}/k_{r0}$ . Such model assumes that holes [ $p_h \propto E_c$ ; the field due to mobile carriers] at the Si/SiON interface tunnel into [probability,  $P_T \sim \exp(\gamma_T E_{ox})$ ] and is captured by Si-H bonds [cross-section,  $\sigma$ ]  $\sim 2\text{\AA}$  away from the interface, leading to a hole-assisted, field-enhanced, thermal generation [rate,  $B \propto \exp(-(E_{F0} - aE_{ox})/k_B T)$ ] of interface traps (Fig. 4). Based on these relations, Eqn. (1) transforms to

$$\Delta N_{IT} = A |Ec|^{2/3} \exp(2\gamma E_{ox}/3) \exp(-nE_{D1}/k_B T) t^n \quad (2)$$

where,  $\gamma = \gamma_T + a/k_B T$ ,  $nE_{D1} = nE_D + 2/3(E_{F0} - E_{R0} - aE_{ox})$ ;  $n \sim 1/6$ .

Eq. (2) suggests the overall activation energy for  $\Delta N_{IT}$  will be

$$E_{IT} = nE_A = nE_D + 2/3(E_{F0} - E_{R0} - aE_{ox}). \quad (3)$$

The assertion that NBTI degradation is field-driven is well known [12, 13] and easily established by Fig. 5 and R-D analysis of wide variety of stress data (Fig. 6). The explicit parameterization of this  $E_{ox}$  dependence is discussed below.

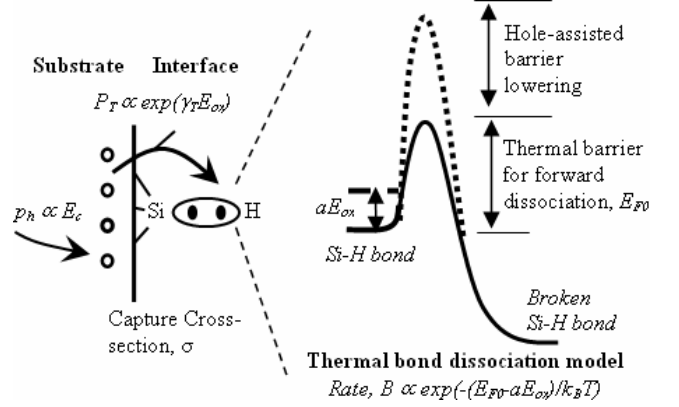


Fig. 4: Hole-assisted, field-enhanced model for Si-H thermal bond dissociation.  $k_{r0}$  is assumed to field independent to first order.

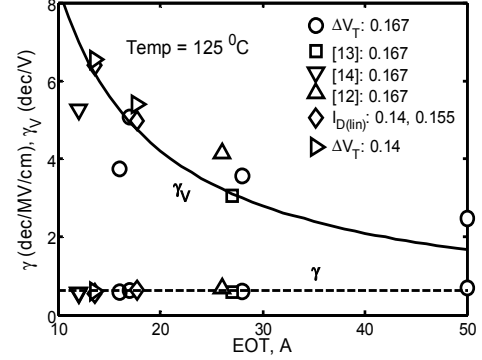


Fig. 5:  $\gamma$  and  $\gamma_V$  (voltage acceleration factor) [12] variation with  $EOT$  ( $n$  is mentioned for each case). For field-dependent NBTI model,  $\gamma$  remains almost constant ( $\sim 0.6 \pm 0.05$ ; dashed line) and  $\gamma_V \sim 1/EOT$  (solid line) [12]. For  $\Delta V_T$  measurements, necessary delay correction is done.

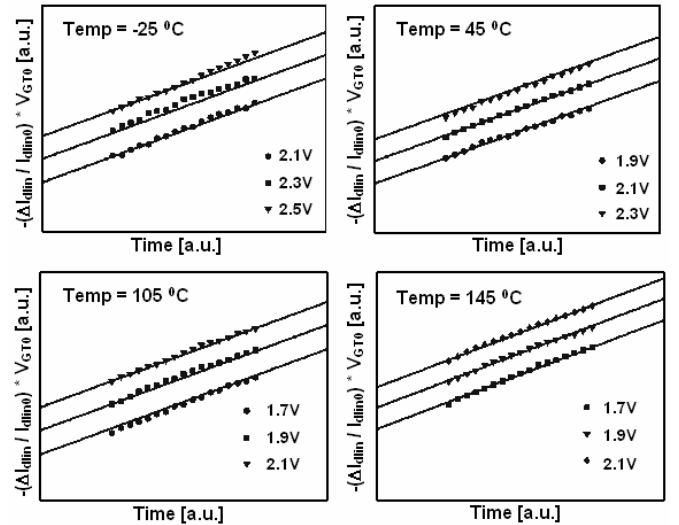


Fig. 6: Fitting  $\Delta V_T$  vs.  $t$  (log-log scale) at different  $T$  and  $V_G$  by (2), validates the universality of the model, which gives  $A, \gamma, n$  as fitting parameters for each temperature (effect of  $E_{D1}$  can be included in  $A$ ). Here  $n = 0.15-0.16$ , indicating  $H_2$  diffusion and  $N_2$  dose: 8%

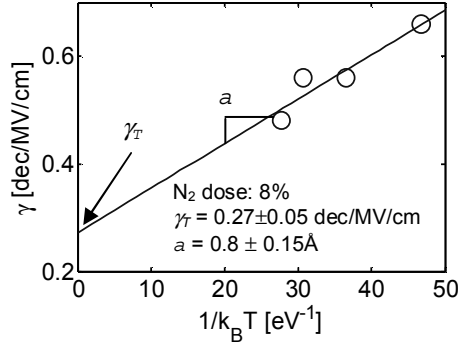


Fig. 7a:  $\gamma$  vs.  $1/k_B T$  plotting.  $\gamma$  obtained from Fig. 6 is used here. The obtained  $\gamma_T = 0.27 \pm 0.05$  dec/MV/cm and  $a = 0.8 \pm 0.15 \text{ \AA}$

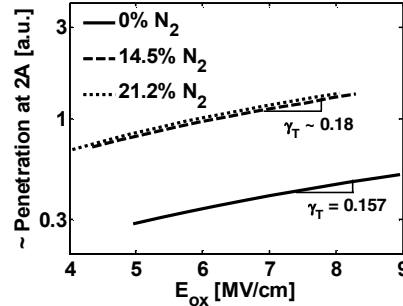


Fig. 7b: Wavefunction penetration at  $2 \text{ \AA}$  (from C-V analysis) modeled using  $A_1 \exp(\gamma_T E_{ox})$ . Analytically [19],  $\gamma_T \sim \sqrt{(m_{ox}/\phi_{th})}$  and  $A_1 \sim \exp \sqrt{(m_{ox}/\phi_{th})}$ , which are consistent with the obtained variations using parameters of Fig. 2.

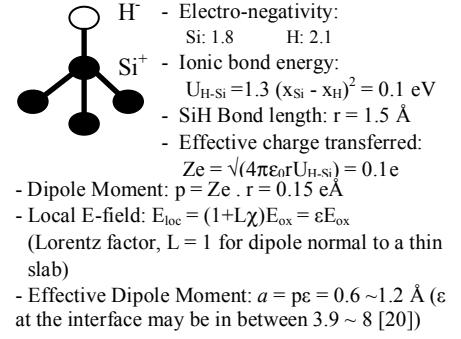


Fig. 7c: Polarization properties of SiH bond, indicating  $a = 0.6 \sim 1.2 \text{ \AA}$ . [21]

### 3.2 Extraction of $\gamma_T$ and $a$

The tunneling coefficient  $\gamma_T$  can either be determined by the *intercept* of measured  $\gamma$  (extracted from Fig. 6) vs.  $1/k_B T$  plot (Fig. 7a) or estimated by the simulation of  $P_T(E_{ox})$  (Fig. 7b, based on parameters from Fig. 2). Both approaches provide comparable results. Similarly effective dipole moment  $a$  can either be measured from *slope* of  $\gamma$  vs.  $1/k_B T$  plot (Fig. 7a) and/or from  $E_A$  vs.  $E_{ox}$  variations (Fig. 9), or be theoretically estimated by molecular modeling approach (Fig. 7c). The summary in Table 1 shows that theory and measurements provide remarkably consistent results. Therefore, part of the  $E_{ox}$ -dependencies of NBTI arises from  $E_{ox}$ -dependence of  $P_T$  and the other part, from  $E_{ox}$ -induced thermal barrier lowering,  $B$  (Fig. 9). This, we believe, is the first consistent interpretation of field-dependence of NBTI degradation within the R-D framework. Previous analysis of delay-contaminated NBTI data [17] suggested experimental values of  $a$  that was 3~4 times larger than theoretically expected values. It appears that similar to resolution of puzzles involving time-exponent ( $n=0.25 \rightarrow 0.16$ ) and temperature dispersion parameter ( $d=0.6 \rightarrow 0$ ) [13], zero-delay measurements and partitioning  $\gamma_T$  into  $\gamma_T + a/k_B T$  (previous work incorrectly assigns  $\gamma = a/k_B T$  [17] or  $\gamma = \gamma_T$  [12] exclusively) hold the key for consistent estimation of field-acceleration of NBTI degradation.

### 3.3 Parametric dependence on %N

The methodologies discussed above have been used to analyze the  $E_{ox}$  and  $T$ -dependent NBTI data at various %N to determine  $\gamma(N) [= \gamma_T(N) + a(N)/k_B T]$  and  $E_A(N)$  (see Fig. 8, 9). While  $\gamma_T$  depends weakly on %N (Fig. 7b),  $a(N) = k_B T [\gamma(N) - \gamma_T(N)] \sim 0.25 dE_A/dE_{ox}$  and  $E_A(N)$  shows systematic variation with %N (Fig. 8, 9 [18]). Hence, our analysis shows that the NBTI degradation with %N variation can be attributed to three factors: reduction of over-all barrier height for dissociation  $nE_{D1}$ , variation in  $A(N)$  and modification of effective dipole moment  $a$ , with negligible contributions from induced hole density ( $p_h$ ).

## 4. Optimization

The parameters of  $\gamma(N)$ ,  $E_A(N)$ ,  $a(N)$ ,  $n(N)$ , derived by fitting the NBTI data and encapsulated within the field-dependent R-D model discussed above, can now be used to calculate  $V_{safe}(N, EOT)$  for any combinations of %N and  $EOT$  (Fig. 10a). Fig. 11 summarizes the results for  $J_G(N, EOT)$  and  $V_{safe}(N, EOT)$ , indicating that co-optimization for DC-NBTI and ITRS-limited  $J_G$  [1] may sometimes require reduced temperature, inclusion of AC effects [23], and NBTI aware design [24]. The diagram can be used to calculate any pair of the variables ( $V_{safe}$ ,  $J_G$ , %N,  $EOT$ ) when the other pair is given. For example, if an IC design requires  $V_{safe} = 1.1 \text{ V}$  and

TABLE 1  
Comparing  $a$  determined by various methodologies

$a$ [ $\text{\AA}$ ]	Obtained from		
0.8	Fit of $\gamma = \gamma_T + a/k_B T$	Fig. 7a	Experiment
0.86 (avg.)	$a = (\gamma - \gamma_T) k_B T$	Fig. 8	
0.98 (avg.)	$a = 0.25 dE_A/dE_{ox}$	Fig. 9	
1.47	$a = (\gamma - \gamma_T) k_B T^*$	[17]	Calculation
0.6~1.2	Si-H polarization	Fig. 7c	
0.65 (avg.)	First Principle†	[22]	

\*  $\gamma$  obtained after delay correction and  $\gamma_T \sim 0.27$ .

† In reference,  $a = 3.2 \text{ \AA}$

‡ For Equilibrium structure

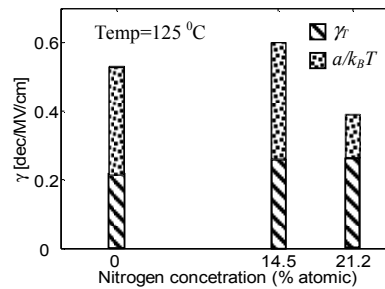


Fig. 8: Variation of  $\gamma$  with %N, after fitting degradation curves for SiON devices, using (2). As stated in the text,  $\gamma$  will have two parts  $\gamma_T \sim 0.225$ ,  $0.27$ ,  $0.27$  (from Fig. 7a and 7b) and  $a/k_B T$ .

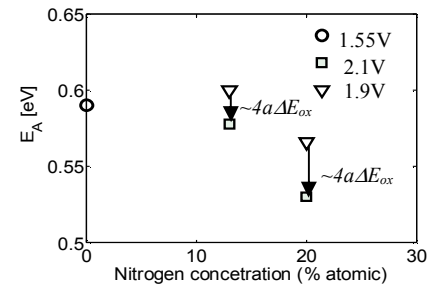


Fig. 9:  $E_A$  determined for various sample. With change in %N, there is a systematic variation in  $E_A$ . There is also slight decrease in  $E_A$  with increase in  $V_G$ , because  $E_A = E_D + 2/3n(E_{F0} - E_{R0} - aE_{ox})$ .

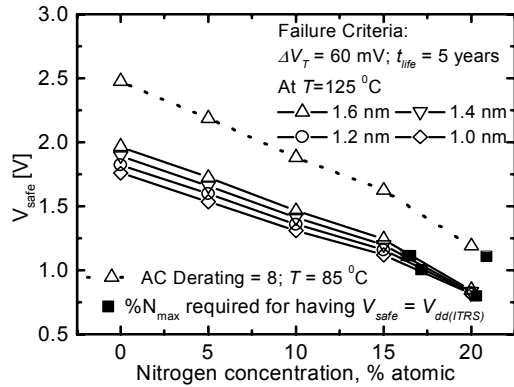


Fig. 10a:  $V_{safe}$  vs. %N at different EOT (Failure criteria:  $\Delta V_T = 60$  mV,  $t_{life} = 5$  years [1]). Dotted line for 1.6 nm EOT indicates the improvement with AC derating = 8 and  $T = 85^\circ\text{C}$ .

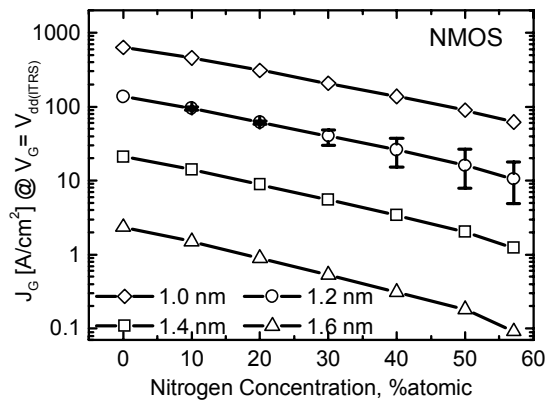


Fig. 10b:  $J_G @ V_{dd(ITRS)}$  for NMOS vs. %N for different EOT. The error bar @ 1.2 nm indicates  $J_G$  variation, expected due to  $\pm 0.5 \text{ \AA}$  error in  $T_{PHY}$ , which is negligible upto 30%  $\text{N}_2$  dose.

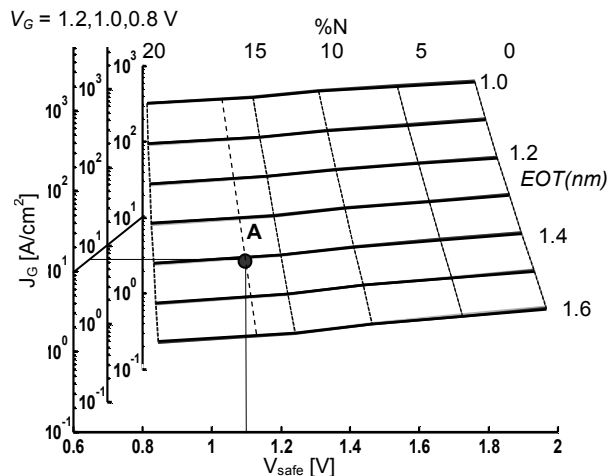


Fig. 11: Design diagram for SiON devices, obtained by using Fig. 10a and 10b ( $J_G @ V_G = 0.8, 1.0, 1.2$  V are considered). Example: choose  $V_G = V_{safe}$  from x-axis and  $J_G$  from  $V_G$  specific y-axis, then read off corresponding %N and EOT. Inclusion of AC effects and temperature reduction will right-shift the diagram, thus enabling co-optimization at smaller EOT.

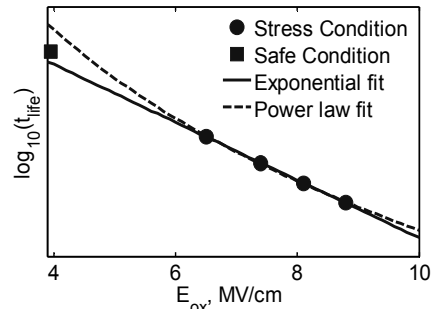


Fig. 12: Comparison between exponential and power law fit of  $\Delta V_T$  data for obtaining safe operating condition using stress data. Procedure in Fig. 3 used for obtaining safe condition. Exponential fit underestimates and power law fit overestimates  $t_{life}$ .

$J_G = 10 \text{ A/cm}^2$  (point A, Fig. 11), the design diagram suggests an EOT = 1.4 nm and %N  $\cong 17\%$ .

## 5. Conclusion

We have analyzed an extensive set of NBTI data to develop the first self-consistent model for field acceleration within R-D framework (compared with other empirical formulation in Fig. 12), constructed a novel design diagram for co-optimization of  $J_G$  and NBTI for arbitrary %N and EOT combinations, and established the DC (AC)-NBTI limited upper limit of %N  $\sim 15\text{-}20\%$  (20-25%) for core CMOS technologies. We anticipate that our analysis would have broad impact for optimization of %N content in sub-2nm gate dielectrics.

## Acknowledgement

This work was done with financial support from TI, AMAT, TSMC, Renesas and SRC. We also acknowledge NCN for providing computational resources.

## References

- [1] ITRS 2005 version.
- [2] M.L. Green et al., *JAP*, 90(5), p. 2057, 2001.
- [3] P.A. Kraus et al., *IEDM*, 52(6), p. 1141, 2005.
- [4] G. Lucovsky et al., *JVST-A*, 18(4), p. 1163, 2000.
- [5] X. Guo et al., *EDL*, 19(6), p. 207, 1998.
- [6] H.Y. Yu et al., *IEDM*, 49(7), p. 1158, 2002.
- [7] S.S. Tan et al., *APL*, 82(12), p. 1881, 2003.
- [8] D.K. Schroder et al., *JAP*, 94(1), p. 1, 2003.
- [9] A. Ghetti et al., *SISPAD*, p. 239, 1999.
- [10] D.M. Brown et al., *JES*, 115(3), p. 311, 1968.
- [11] J.R. Shallenberger et al., *JVST-A*, 17(4), p. 1086, 1999.
- [12] M.A. Alam et al., *Micro. Reliability*, 45(1), p. 71, 2005.
- [13] D. Varghese et al., *IEDM*, p. 684, 2005.
- [14] S. Chakravarthi et al., *IRPS*, p. 273, 2004.
- [15] A.T. Krishnan et al., *IEDM*, p. 688, 2005.
- [16] C.L. Chen et al., *IRPS*, p. 704, 2005.
- [17] K.O. Jeppson et al., *JAP*, 48(5), p. 2004, 1977.
- [18] J.P. Campbell et al., *IRPS*, p. 442, 2006.
- [19] Register et al., *APL*, 74(3), p. 457, 1999.
- [20] Muller et al. *Nature*, 399(6738), 1999.
- [21] Mcpherson et al., *JAP*, 84(3), p. 1513, 1998.
- [22] Ushio et al., *JAP*, 97(8), 2005.
- [23] M.A. Alam, *IEDM*, p. 345, 2003.
- [24] B.C. Paul et al., *EDL*, 26(8), p. 560, 2005.